Complex near-surface anomalies are one of the main onshore challenges facing seismic data processors. Refraction tomography is becoming a common technology to estimate an accurate near-surface velocity model. This process involves picking the first arrivals of refracted waves. One of the main challenges with refraction tomography is the low signal-to-noise ratio characterizing the first-break waveform arrivals, especially for the far-offset receivers. This is especially evident in data recorded using reflection acquisition geometry. This low signal-to-noise ratio is caused by signal attenuation due to geometrical spreading of the seismic wavefield, near-surface-generated noise, and amplitude absorption. Super-virtual refraction interferometry improves the quality of the first-break picks by enhancing the amplitude of the refracted waves and attenuating the amplitude of the random noise.

The theory of refraction interferometry was developed by Dong et al. (2006) and later Bharadwaj and Schuster (2010) successfully applied the technique. Dong et al. correlate a pair of traces to give $\phi(A|B)$, where A and B are the geophone positions and $x$ is the source position. The resulting virtual trace has a virtual refraction arrival time of $\tau_{cv} - \tau_{cv'}$. Repeating this procedure for any postcritical source position leads to a virtual trace of the same virtual refraction traveltime. Stacking the correlated traces over all postcritical source positions yields a trace with a virtual refraction event that has an enhanced signal-to-noise ratio. This enhancement can be as high as $\sqrt{N}$, where $N$ is the number of sources that contribute to the generation of this particular virtual head wave. They demonstrate this method on land data shot over a salt dome in Utah and later Nichols et al. (2010) showed its effectiveness in a hydrogeophysical research site in Idaho.

A problem with refraction interferometry is that, if only the head wave arrivals are correlated with one another, the virtual head-wave trace has the correct moveout pattern. It has an unknown excitation time, so as a remedy, Dong et al. suggested that the source be “virtually” relocated to the surface by calibrating the virtual stacked refraction trace to an observed travelttime in the raw data. Another problem is that correlation of traces typically decreases the source-receiver offset of the virtual trace because traveltimes are subtracted and are associated with shorter raypaths. To overcome this, Bharadwaj and Schuster, Mallinson et al. (2011), and Hanafy et al. (2011) presented an extension of refraction interferometry so that the receiver spread could be extended to its maximum recording extent, and the absolute arrival time can be properly accounted for. This new method creates virtual far-offset refraction arrivals by

Figure 1. (a) Cross-correlating the traces of refracted arrivals at R1 and R2 would redatum the source to the refractor at R1'. This virtual source will have a trace traveling from R1' to R2 with a negative excitation time equal to the time between R1 and R1'. (b) Cross-correlating the traces using a different source position (S2) but with the same receiver pair (R1 and R2) generates the same virtual source trace as in (a). (c) As long as the receiver pair (R1 and R2) is present in a particular shot gather, then it is possible to obtain a trace that has the same raypath as in (a) for stacking.
same receiver pair. Figures 1b and 1c show that the same virtual source trace can be obtained using different shot positions with the same receiver pair. Hence the traces generated using the virtual sources are stacked to increase the signal-to-noise ratio. The trace of the virtual source and an actual trace that travels to $R_3$ are convolved to redatum the source back to the surface, remove the effect of negative time, and achieve a second stack. The reciprocity equation of convolution type in the far-field approximation gives the new redatumed Green's function according to the following expression:

$$G(R_4|S)_{\text{super}} = 2i k_{\text{virtual}} \int G(R_3'|S)_{\text{virtual}} G(R_4'|S)_{\text{virtual}} dS$$  \hspace{1cm} (2)$$

where $i = \sqrt{-1}$, $k$ is the wavenumber of the head wave, and $G(R_3|S)$ and $G(R_4|S)$ are the recorded Green's functions from source $S$ to receivers $R_3$ and $R_4$, respectively. The result of the cross-correlation $G(R_4|R_3')_{\text{virtual}}$ is a virtual source located at $R_3'$ (the point where the raypaths to $R_3$ and $R_4$ diverge). The integration surface $(S_0 + S_{\infty})$ reduces to $S_0$ due to the Wapenaar anti-radiation condition which states that for a sufficiently heterogeneous medium which assumes little interactions at infinity, the recorded wavefield is negligible.

The excitation time of the virtual source is equal to the time it takes a seismic wave to travel between $R_3$ and $R_3'$ (Mallinson et al.). Figure 1a shows the result of cross-correlation where the dashed line indicates the negative time component. The virtual source is now independent of source position; thus the virtual source can be generated for different sources that have the same receiver pair. Figures 1b and 1c show that the same virtual source trace can be obtained using different shot positions with the same receiver pair. Hence the traces generated using the virtual sources are stacked to increase the signal-to-noise ratio. The trace of the virtual source and an actual trace that travels to $R_3$ are convolved to redatum the source back to the surface, remove the effect of negative time, and achieve a second stack. The reciprocity equation of convolution type in the far-field approximation gives the new redatumed Green's function according to the following expression:

$$G(R_4|S)_{\text{super}} = 2i k_{\text{virtual}} \int G(R_3'|S)_{\text{virtual}} G(R_4'|S)_{\text{virtual}} dS$$  \hspace{1cm} (2)$$

where $i = \sqrt{-1}$, $k$ is the wavenumber of the head wave, and $G(R_4'|S)_{\text{virtual}}$ is the stacked trace obtained by Equation 1. The superscript in the new redatumed Green's function $G(R_4'|S)_{\text{super}}$ is to indicate that this trace is different from the original recorded Green's function $G(R_4'|S)$, which is used in the reciprocity equation of correlation type (Equation 1). Figure 2a shows the redatuming process occurring in the convolution phase. Notice that, just as in the cross-correlation phase, the redatumed event in the convolution phase is now independent of virtual source position. Therefore it is possible to stack over the virtual source position. This is the second stack applied to the refracted data in order to achieve superior signal-to-noise ratio. The trace of the virtual source and an actual trace that travels to $R_3$ are convolved to redatum the source back to the surface, remove the effect of negative time, and achieve a second stack. The reciprocity equation of convolution type in the far-field approximation gives the new redatumed Green's function according to the following expression:

$$G(R_4|S)_{\text{super}} = 2i k_{\text{virtual}} \int G(R_3'|S)_{\text{virtual}} G(R_4'|S)_{\text{virtual}} dS$$  \hspace{1cm} (2)$$

where $i = \sqrt{-1}$, $k$ is the wavenumber of the head wave, and $G(R_4'|S)_{\text{virtual}}$ is the stacked trace obtained by Equation 1. The superscript in the new redatumed Green's function $G(R_4'|S)_{\text{super}}$ is to indicate that this trace is different from the original recorded Green's function $G(R_4'|S)$, which is used in the reciprocity equation of correlation type (Equation 1). Figure 2a shows the redatuming process occurring in the convolution phase. Notice that, just as in the cross-correlation phase, the redatumed event in the convolution phase is now independent of virtual source position. Therefore it is possible to stack over the virtual source position. This is the second stack applied to the refracted data in order to achieve superior signal-to-noise ratio. The trace of the virtual source and an actual trace that travels to $R_3$ are convolved to redatum the source back to the surface, remove the effect of negative time, and achieve a second stack. The reciprocity equation of convolution type in the far-field approximation gives the new redatumed Green's function according to the following expression:

$$G(R_4|S)_{\text{super}} = 2i k_{\text{virtual}} \int G(R_3'|S)_{\text{virtual}} G(R_4'|S)_{\text{virtual}} dS$$  \hspace{1cm} (2)$$

where $i = \sqrt{-1}$, $k$ is the wavenumber of the head wave, and $G(R_4'|S)_{\text{virtual}}$ is the stacked trace obtained by Equation 1. The superscript in the new redatumed Green's function $G(R_4'|S)_{\text{super}}$ is to indicate that this trace is different from the original recorded Green's function $G(R_4'|S)$, which is used in the reciprocity equation of correlation type (Equation 1). Figure 2a shows the redatuming process occurring in the convolution phase. Notice that, just as in the cross-correlation phase, the redatumed event in the convolution phase is now independent of virtual source position. Therefore it is possible to stack over the virtual source position. This is the second stack applied to the refracted data in order to achieve superior signal-to-noise ratio.
volves two stacking operations: one in the correlation phase and one in the convolution phase.

**Synthetic data example**

Prior to applying the super-virtual refraction interferometry theory on real data acquired using reflection geometry, the technique was tested on synthetic data with similar geometry to the real data. The synthetic data are generated using a finite-difference solution to the acoustic wave equation. Moreover, the number of sources generated was 23 with a source spacing of 10 m; each shot gather contained 243 receivers with 10-m interval spacing. Figure 3a shows the results obtained using a finite-difference modeling (note the green plus signs indicate the first picks). Notice that the refraction jumps are clearly visible. Refraction jumps refer to the crossover distance where a different refractor arrival becomes the first arrival.

Random noise was added to all the synthetic traces so that the refracted arrivals are no longer clearly visible. Figure 3b shows the same common-shot gather as in Figure 3a but with random noise added. It is difficult to follow the refracted first arrival and the refraction jump is no longer visible. This is because at this offset the noise level is becoming comparable to the signal level which is making the process of distinguishing between the signal and noise difficult. The mean value of signal-to-noise ratio of the common-shot gathers was calculated to be two.

Super-virtual refraction interferometry was applied to the synthetic data. This increased the signal-to-noise ratio dramatically due to the two stacking operations following each redatuming step. Therefore, the previously masked first-break refracted arrivals and the refraction jump were now clearly visible and easier to pick. Figure 3c shows the result of applying super-virtual refraction interferometry on the same traces as in Figures 3a and 3b.

One can compare the traveltime picks versus offset for all three data sets for one common-shot gather (in our case we chose CSG 12). In the noise-free synthetic data set, it was possible to pick all the first arrivals (up to 2430 m). After adding
noise, the first-break picking could be reliably performed up to an offset of about 480 m. After applying super-virtual refraction interferometry, picking could be done up to an offset of 910 m. Thus, super-virtual refraction interferometry was able to extend the first-break picking to a usable offset range by almost 90%. Figure 4a shows the traveltimes of both the super-virtual and noise-free synthetic picks. All the picks fall within a quarter of a cycle (T/4). This could be considered as a quality assurance because the traveltime picks are almost identical to the picks obtained from a common-shot gather with pure signal. Figure 4b shows a histogram of the difference between the picks. Notice that almost all the events have a traveltime difference of less than a quarter of a cycle (red line indicates the quarter of the cycle).

**Field data example**

Super-virtual refraction interferometry has been tested on a land data set obtained in Saudi Arabia. The geometry is designed for reflected-wave acquisition. As in the synthetic example, the number of receivers is 243 with 10-m interval spacing; and the shot-interval spacing is 10 m. Due to irregularities in the source line position, only 23 common-shot gathers are chosen to apply the methodology.

A window is applied around the region of the first arrivals. The smaller the window, the less the artifacts appear after applying super-virtual refraction interferometry (Figure 5a). The first arrivals of the traces are clear up to 600-m offset; beyond this offset, the signal-to-noise ratio is very low. At approximately 750 m, it is difficult to distinguish the signal from the noise; thus picking the first arrival is not possible using conventional methods. Figure 5a illustrates a zoomed version of the common-shot gather.

Super-virtual refraction interferometry is applied to the windowed land data set. The signal-to-noise ratio of the refracted arrivals is now significantly increased and the random noise is attenuated. This enables picking the first arrival at farther offsets. Figure 5b shows the results after applying super-virtual refraction interferometry. To validate the accuracy of the picked traveltimes, first-arrival times are picked in the super-virtual gather and compared to the raw data picks (Figure 6a). The difference in the common traveltimes is mostly within a quarter of a period (T/4).

The first-arrival picks are inverted to obtain a near-surface 2D velocity tomogram. The obtained tomograms show that...
Near-surface measurements in exploration geophysics

The aperture has almost doubled and the reliability increased (due to an increased ray illumination). Also, because the events of interest are dipping, the depth of investigation has also increased after adding longer-offset traveltime picks. Generally, there are more features in the super-virtual section due to the extra picks (Figure 7). It should be noted that cross-correlation creates “ringier” super-virtual seismograms, which may be an additional complication but that can be addressed with an additional deconvolution step.

Wavelet distortion

When applying the cross-correlation and convolution processes, the seismic wavelet is no longer preserved. What is obtained is the result of cross-correlating and convolving two similar wavelets. This distorts the wavelet and creates side lobes that may confuse the first-break picker (Figure 8). One method to overcome this obstacle is to pick the first arrival for a couple of near-offset traces from the raw section (i.e., before applying super-virtual refraction interferometry). Then apply super-virtual refraction interferometry on the data set, and load the previous picks on top of the super-virtual section. Subsequently the refracted arrival will be identified and picking can resume to longer offsets. Another method to overcome this effect is to add an additional deconvolution step.

Extending from 2D to 3D

Extending super-virtual refraction interferometry from 2D to 3D is a great challenge. The problem is that current 3D acquisition surveys have few sources that generate common-refractor raypaths between the receivers due to the sparseness of source-receiver field configuration. Alternatively, in a dense 3D survey, there is a higher probability of common refractor overlap and thus super-virtual refraction interferometry is expected to enhance the first breaks as well.

Conclusion

Using the super-virtual refraction interferometry method, the signal-to-noise ratio of far-offset refracted wave arrivals is expected to increase by approximately a factor of $\sqrt{N}$ where $N$ is the number of actual sources and virtual sources used in the interferometric summation. Super-virtual refracted arrivals are generated in two steps; the first step involves the cross-correlation and stacking of the data to generate traces with virtual head-wave arrivals and the second step entails the convolution and stacking of the data with the virtual traces to enhance head-wave arrivals.

Super-virtual refraction interferometry has been applied to a data set with a complex near surface from Saudi Arabia. Although the acquisition geometry employed is optimized for seismic reflection acquisition, the method has proven useful to extract and enhance the refracted arrivals in the data set. The resulting tomograms show that the aperture and depth of investigation have increased, producing more reliable near-surface models.

References


Acknowledgments: We thank Saudi Aramco for the opportunity to work on this project and publish the results; in particular we thank Panos Kelamis and Tim Keho for their support of this project. We are grateful to Andrey Bakulin, Constantine Tsingas, and Emad Hemyari (all from Saudi Aramco) for their valuable comments and suggestions, which improved the quality of this article. We are thankful to Gerard Schuster (KAUST) for introducing us to this technology and helping us throughout the project.

Corresponding author: abdulrahman.shubail@aramco.com