High optical and structural quality of GaN epilayers grown on $(2 \times 01) \beta\text{-Ga}_2\text{O}_3$

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High optical and structural quality of GaN epilayers grown on $(\bar{2}01)$ β -Ga₂O₃

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Producing highly efficient GaN-based optoelectronic devices has been a challenge for a long time due to the large lattice mismatch between III-nitride materials and the most common substrates, which causes a high density of threading dislocations. Therefore, it is essential to obtain alternative substrates with small lattice mismatches, appropriate structural, thermal and electrical properties, and a competitive price. Our results show that $(\bar{2}01)$ oriented β -Ga₂O₃ has the potential to be used as a transparent and conductive substrate for GaN-growth. Photoluminescence spectra of thick GaN layers grown on $(\bar{2}01)$ oriented β -Ga₂O₃ are found to be dominated by intense bandedge emission. Atomic force microscopy studies show a modest threading dislocation density of $\sim 10^8 \text{ cm}^{-2}$. X-ray diffraction studies show the high quality of the single-phase wurtzite GaN thin film on $(\bar{2}01)$ β -Ga₂O₃ with in-plane epitaxial orientation relationships between the β -Ga₂O₃ and the GaN thin film defined by $(010) \beta\text{-Ga}_2\text{O}_3 \parallel (11\bar{2}0) \text{ GaN}$ and $(\bar{2}01) \beta\text{-Ga}_2\text{O}_3 \parallel (0001) \text{ GaN}$ leading to a lattice mismatch of $\sim 4.7\%$. Complementary Raman spectroscopy indicates that the quality of the GaN epilayer is high. © 2014 AIP Publishing LLC. [<http://dx.doi.org/10.1063/1.4891761>]

III-nitride materials, in particular, GaN, are attractive because of their large direct energy bandgap and their high thermal stability properties. They are used in numerous optoelectronic and electronic applications, such as laser diodes (LDs), visible and ultraviolet (UV) detectors, and light-emitting diodes (LEDs).^{1–4} Despite the wide use of GaN materials, much effort is still being devoted to improving the optical efficiency and stability of III-nitride materials. A major issue concerning device performance of III-nitrides is the introduction of high densities of threading dislocations (TDs) and other defects such as stacking faults, grain boundaries, and point defects that act as nonradiative recombination “deep state defects.”^{5–7} These deep sites trap carriers and significantly reduce radiative electron-hole pair recombinations.⁸

III-nitride materials are fabricated on sapphire (Al₂O₃), SiC, or Si substrates. Researchers have achieved improvements in growth techniques, including dislocation reduction due to epitaxial lateral overgrowth or even growth on bulk GaN.⁹ However, sapphire, the most commonly used transparent substrate is electrically insulating and has a 14% lattice mismatch, which leads to high dislocation densities (TDDs).¹⁰ Another routinely used substrate, SiC, has a much smaller lattice mismatch (3.1%)¹¹ due to higher thermal and electrical conductivities, providing good stability. However, the poor surface morphology of SiC creates defects in the GaN epilayer. In addition, SiC is expensive and has high optical losses due to the lack of transparency.¹² It is therefore important to find alternative suitable substrates that have good lattice-match to provide high-quality materials with low density of TDs and point defects and appropriate

thermal, optical, and electrical properties with simplified device structure.

Monoclinic β -Ga₂O₃ is a good candidate as a substrate for the growth of III-nitride materials, in particular, GaN. The lattice mismatch between β -Ga₂O₃ and GaN was found to be small (a minimum in-plane mismatch of about 2.6% for the in-plane epitaxial relationship $(011) \beta\text{-Ga}_2\text{O}_3 \parallel (10\bar{1}0) \text{ GaN}$).¹³ Ga₂O₃ has other advantages over other substrates, including high conductivity (Sn dopants increases its conductivity to be $\sim 1 \text{ S cm}^{-1}$)¹⁴ and high transparency in the visible and UV spectral regions due to its wide band gap (4.8 eV),¹⁵ thus combining the advantages of Al₂O₃ and SiC substrates. Preliminary studies on the molecular beam epitaxial (MBE) growth of GaN on monoclinic (100) β -Ga₂O₃ have already been reported, although they showed low GaN quality and low reproducibility¹⁶ because of the cleavage of the Ga₂O₃ (100) plane, which causes the striping of the epilayer and substrate. Currently, the growth of GaN on the $(\bar{2}01)$ plane of Ga₂O₃ is under investigation for the development of InGaN LEDs.¹⁷

In this work, we study a high-quality GaN epilayer grown on monoclinic β -Ga₂O₃ by metal organic chemical vapor deposition (MOCVD). We show that $(\bar{2}01)$ β -Ga₂O₃ is a promising substrate for III-nitride-based devices as we found that the in-plane lattice mismatch between the $(\bar{2}01)$ plane of β -Ga₂O₃ and the (0002) plane of GaN is as low as 4.7%, which results in a high-quality GaN epilayer with a low TDD and a high internal quantum efficiency.

A 0.48 μm thick AlN buffer layer was grown on $(\bar{2}01)$ -oriented β -Ga₂O₃ at 500 °C under a N₂ and NH₃ atmosphere. Then, after changing the carrier gas from N₂ to H₂, the substrate temperature was increased to 1020 °C to grow a layer of GaN followed by another increase to 1080 °C to grow

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another GaN layer. The total thickness of GaN on the AlN buffer layer was $3\ \mu\text{m}$. The MOCVD growth conditions are detailed in Ref. 18. A commercial (Lumilog) GaN film with a low TDD ($4 \times 10^7\ \text{cm}^{-2}$) grown on sapphire (GaN/Al₂O₃) by MOCVD was used to compare the quality of the sample.

X-ray diffraction (XRD) measurements were carried out to examine the epitaxial relationship of the film and substrate, the structural properties, and the strain. The XRD measurements were performed on a Bruker D8 AXS diffractometer using a Cu K_{α1} line, a 2-bounce Ge (220) monochromator, and a Göbel mirror. Asymmetric reciprocal space maps (RSMs) were acquired using a $0.1\ \mu\text{m}$ wide slit placed in front of a scintillation detector. Rocking curves (RCs) and pole figures were acquired using the open detector. Room temperature (RT) Raman measurements were carried out in the Horiba LabRam ARAMIS micro-Raman system with backscattering geometry. The samples were excited by a 473 nm diode laser and the data were collected by a monochromator equipped with 1800 lines/mm grating and a CCD camera cooled by liquid nitrogen. The surface morphology, roughness, and microstructural defects of the epilayers were investigated by atomic force microscopy (AFM) using the scanning probe microscope, Agilent 5400. Photoluminescence (PL) was measured to investigate the optical properties of the GaN film using a 325 nm He-Cd laser. The spectra were detected by an Andor monochromator attached to a CCD camera. The samples were mounted in a closed-cycle Helium cryostat for low-temperature PL (6 K).

The 2θ - θ XRD scan in Fig. 1(a) shows the diffraction pattern of the GaN/ β -Ga₂O₃ sample, revealing the β -Ga₂O₃ substrate diffraction peaks at (201), (402), (603), (804), (1005). Symmetric reflections (0002), (0004) and (0006) were observed, as expected for a single crystalline, c-plane oriented wurtzite GaN film. The GaN lattice parameters of the two samples were determined using the Bond method¹⁹ by measuring the (0004) and (10 $\bar{1}$ 4) symmetric and asymmetric RCs. The measurements yielded values of $c=5.187(1)\ \text{\AA}$ and $a=3.185(1)\ \text{\AA}$ and $c=5.192(1)\ \text{\AA}$ and $a=3.181(1)\ \text{\AA}$ for GaN/Ga₂O₃ and GaN/Al₂O₃, respectively, showing slight

compressive strain in both samples ($\epsilon^{\text{xx}} = -0.12(3)\%$ and $\epsilon^{\text{zz}} = 0.04(2)\%$ for GaN/Ga₂O₃ and $\epsilon^{\text{xx}} = -0.24(3)\%$ and $\epsilon^{\text{zz}} = 0.14(2)\%$ for GaN/Al₂O₃). The strain was calculated assuming strain-free lattice parameters of $a=0.18878\ \text{\AA}$ and $c=5.18500\ \text{\AA}$.²⁰ The strain ratio, $\epsilon^{\text{zz}}/\epsilon^{\text{xx}}$, should be constant in samples where only biaxial strain is present and for this case a value of $\epsilon^{\text{zz}}/\epsilon^{\text{xx}} = -0.6$ was estimated by Kisielowski *et al.*²⁰ In the present case, we find strain ratios of $\epsilon^{\text{zz}}/\epsilon^{\text{xx}} = -0.3(1)$ and $\epsilon^{\text{zz}}/\epsilon^{\text{xx}} = -0.6(1)$ for GaN/Ga₂O₃ and GaN/Al₂O₃, respectively, suggesting that a hydrostatic strain component cannot be disregarded in our GaN/Ga₂O₃ sample. Such hydrostatic strain may be introduced by defects or impurities. Figs. 1(b)–1(d) show the XRD RCs around the GaN (0002) and (10 $\bar{1}$ 5) reflections in the samples grown on Ga₂O₃ and Al₂O₃. The FWHM for the plots are shown in the brackets. For (10 $\bar{1}$ 5), the measurements were performed with asymmetric as well as skew symmetric geometries. The values for the integral breadth ranging from 0.10° to 0.14° for GaN/Ga₂O₃ are indicated in the figures and reveal the good quality of both samples with typical values for state-of-the-art films. The broadening of RCs in GaN is typically associated with the mosaicity and dislocation density of the films. In particular, the (0002) RCs are broadened by mosaic tilt caused by screw-type and mixed dislocations (as well as the lateral coherence length), while the (10 $\bar{1}$ 5) reflection is also sensitive to edge-type dislocations (causing twist) when measured with the skew symmetric geometry.²¹ The two samples show very similar behavior with only slightly higher values for the sample grown on Ga₂O₃ suggesting slightly higher dislocation densities, an impressive result for initial growth experiments with little optimization of growth parameters. The mosaicity was further studied by using the RSM of the (10 $\bar{1}$ 5) asymmetric reflection as shown in Figs. 1(e) and 1(f). From the broadening of the reciprocal lattice points, the tilt and lateral coherence length, $L_{||}$, were estimated using the geometrical method described by Fewster.²² In both samples, the tilt was very low, suggesting that the lateral coherence length contributes considerably to the broadening (see Figs. 1(e) and 1(f)). In fact, the tilt of the GaN/Ga₂O₃ layer was lower than that of the GaN/Al₂O₃ reference sample while the lateral coherence length was smaller, which also suggests that the main broadening of the (0002) RC (Fig. 1(b)) in GaN/Ga₂O₃ was due to a small lateral coherence length. XRD pole figures were acquired on the GaN/Ga₂O₃ sample by fixing the detector at an angle of $2\theta = 48.27^\circ$ to measure the reflections of the {10 $\bar{1}$ 2} family of planes for GaN and the {5 $\bar{1}$ 1} and {5 $\bar{1}$ 0} planes for the case of β -Ga₂O₃ simultaneously (Fig. 2). The epitaxial orientation relationships between the β -Ga₂O₃ and the GaN thin film are defined by (010) β -Ga₂O₃ || (11–20) GaN and (201) β -Ga₂O₃ || (0001) GaN, resulting in a lattice mismatch of $\sim 4.7\%$.

The termination of TDs in the epilayer surface is indicated by small pits in the AFM scan. A wide scan of the surface morphology is shown in Fig. 3(a). Fig. 3(b) shows a higher-resolution AFM image. The average density of TDs is found to be $\sim 10^8\ \text{cm}^{-2}$ (by averaging over 20 images taken at different positions on a 2-in. wafer), which is relatively low compared with the TDDs reported in early works on GaN epilayers grown on other substrates.^{23,24} The TDD of GaN epilayers grown on Al₂O₃ varies typically in the range

FIG. 1. (a) The θ - 2θ XRD curve of the GaN/Ga₂O₃ epilayer. (b) The rocking curve around (0002) for GaN/Ga₂O₃ and GaN/Al₂O₃ epilayers and FWHM is shown in the bracket. (10 $\bar{1}$ 5) Rocking curves taken with asymmetric and skew symmetric geometries for (c) GaN/Ga₂O₃ and (d) GaN/Al₂O₃ having a FWHM as shown in the brackets. (10 $\bar{1}$ 5) Reciprocal space maps for (e) GaN/Ga₂O₃ and (f) GaN/Al₂O₃.

FIG. 2. A pole figure showing the $\{10\bar{1}2\}$ family of planes for GaN and the $\{5\bar{1}1\}$ and $\{5\bar{1}0\}$ planes for the case of Ga_2O_3 .

of 10^8 – 10^{10} cm^{-2} and our reference sample shows a low TDD of $4 \times 10^7 \text{ cm}^{-2}$. This low value is in agreement with the slightly lower RC widths described above compared to the GaN/ Ga_2O_3 layer. More sophisticated research approaches like patterning the substrates and optimized multistep MOCVD techniques have reduced the TDD to a value of $\sim 10^{-6} \text{ cm}^{-2}$.^{25,26} We suggest that this TDD can be reduced significantly using $(\bar{2}01)$ β - Ga_2O_3 substrates and controlling the growth process and specifications of the buffer layer and the film.

Large differences in the magnitude of Burgers vectors, especially between edge-type and screw/mixed-type dislocations, imply that the size of the surface pits should be different depending on the type of dislocation.²⁷ We could not clearly identify any screw type (the largest pits) in agreement with the low tilt values estimated from XRD. We mainly noticed mixed dislocations (medium pits) and edge dislocations (smallest pits) at a ratio of 2:3 per μm^2 . Medium-sized pits have a maximum depth of $\sim 7 \text{ nm}$ and a width of $\sim 40 \text{ nm}$; small pits are about 20 nm wide and have a maximum depth of $\sim 0.6 \text{ nm}$. The stepped surface morphology exhibits atomic terraces $\sim 0.2 \text{ nm}$ in height. The root mean square (RMS) roughness is found to be as low as $\sim 0.68 \text{ nm}$ over $100 \mu\text{m}^2$, which indicates a very smooth surface compared to the reference sample (with a RMS roughness of $\sim 0.87 \text{ nm}$).

Figs. 4(a) and 4(b) show RT and low-temperature (6 K) PL spectra of GaN/ Ga_2O_3 , respectively, compared to GaN/ Al_2O_3 . The RT PL spectrum of GaN/ β - Ga_2O_3 has a dominant intensity near the bandedge emission centered at 3.41 eV and a very weak yellow luminescence emission. The bandedge emission presents a combination of free exciton transition, donor-bound exciton transition, and donor-to-valence band transition.^{28–30} Surprisingly, the ratio of the PL intensity at 6 K to that at RT, found to be $\sim 2.7:1$, is better than that of the GaN/ Al_2O_3 ($\sim 20:1$) and is similar to that in undoped GaN/ Al_2O_3 epilayers

FIG. 3. (a) This $5 \times 5 \mu\text{m}^2$ AFM image of the GaN/ Ga_2O_3 and (b) $1 \times 1 \mu\text{m}^2$ AFM image of the GaN/ Ga_2O_3 sample.

FIG. 4. (a) (Log scale) PL spectra at RT and (b) (linear scale) normalized PL spectra at 6 K, and the absorption spectra (unnormalized spectra for reflection) (c) at RT for both the GaN/ Ga_2O_3 and GaN/ Al_2O_3 epilayers (the inset shows the absorption spectrum of GaN/ Ga_2O_3 at 6 K).

reported in the literature.³¹ The bandedge intensity of GaN/ Ga_2O_3 wafer is 1.5 times stronger at RT than that from the GaN/ Al_2O_3 wafer, although the TDD of the later is lower. This suggests that other defects that create a high density of nonradiative deep states (such as point defects or other structural defects) are suppressed significantly, resulting in much higher intensity. Fig. 4(b) shows an intense band edge emission at 6 K. The bandedge peak of GaN/ Ga_2O_3 is redshifted compared to that of the film on Al_2O_3 , in agreement with the lower compressive strain observed by XRD in the former. However, the peak wavelength can also be influenced by doping or defects.²⁰ The inset of Fig. 4(b) shows donor acceptor pair (DAP) peaks at 6 K within the energy range from 3.28 eV to 3 eV . A series of peaks are observed, with the first peak centered at 3.29 eV followed by the longitudinal optical (LO)-phonon replicates at low energies, with an energy difference of about $\sim 88 \text{ meV}$. In general, we observed that GaN/ β - Ga_2O_3 is mechanically rigid and the PL intensity remains the same even in cut pieces compared with the GaN/ Al_2O_3 .

The optical absorption spectra of GaN/ Ga_2O_3 and GaN/ Al_2O_3 at RT are shown in Fig. 4(c). The band gap of GaN/ Ga_2O_3 epilayer is measured as $\sim 3.4 \text{ eV}$ at RT, slightly blue-shifted compared to bulk GaN (3.39 eV at RT).³² This slight blueshift can be due to the compressive stress. The near UV absorption is higher for the GaN epilayer grown on Ga_2O_3 than that grown on Al_2O_3 , which is due to the absorption from Ga_2O_3 (figure not shown). A slight tail is shown in the absorption spectra of GaN/ Ga_2O_3 at RT and 6 K (inset of Fig. 4(c)), compared to that of GaN/ Al_2O_3 . The exponential behavior in the tail of the absorption edge can be attributed to structural disorder, point defects, excitonic transitions, or inhomogeneous strain in the film,³³ which can cause a broadening of the PL bandedge peak.

Fig. 5(a) shows the Raman spectrum of the GaN/ Ga_2O_3 epilayer and that of the β - Ga_2O_3 substrate as a reference to identify the GaN peaks (Fig. 5(b)). The Raman peaks of β - Ga_2O_3 substrate (at 169.26 , 199.81 , 319.7 , 346.89 , 416.04 , and 659.95 cm^{-1}) are marked by asterisks. It is known that

FIG. 5. The RT Raman spectra of (a) the GaN/Ga₂O₃ and (b) bulk β -Ga₂O₃ substrate.

β -Ga₂O₃ is a monoclinic structure with space group symmetry, $C_{2/m}$.³⁴ The results for (201)-oriented β -Ga₂O₃ agree with the results observed by Dohy *et al.*³⁵ As shown in Fig. 5, not all the optical modes of β -Ga₂O₃ are observed.³⁴ However, our observation shows that all the translational modes of bulk β -Ga₂O₃ are redshifted in the GaN/Ga₂O₃ epilayer by 1.5–2 cm⁻¹ when compared to bulk β -Ga₂O₃. Wurtzite GaN is a tetrahedrally coordinated semiconductor with the space group of C_{6v}^4 and C_{3v} symmetry. It is expected that Raman active modes (one A_1 , one E_1 , and two E_2) would be observed in the GaN epilayer.³⁶ The modes observed in this epilayer are E_2 (high) and E_2 (low) modes. The E_2 (high) mode peak is clearly observed (at 569.04 cm⁻¹), as shown in Fig. 5(a). The symmetry E_2 (low) mode at 143.2 cm⁻¹ is observed and overlaps partially with the 144.4 cm⁻¹ substrate mode (Ga₂O₃). The detected E_2 modes verify the backscattering through the c-axis of wurtzite GaN.³⁷ We observe that the positions of these peaks compared with the positions of the peaks for GaN grown on Al₂O₃ are dependent on the substrate due to different strain states for different samples (data not shown). The E_2 (high) mode has been demonstrated to be sensitive to biaxial strain in GaN epilayers.³⁷ Here, we can predict that there is a slight compressive strain because the peak exhibits a blueshift of 1.04 cm⁻¹ compared with that of bulk GaN (568 cm⁻¹).³⁷ In contrast, the Raman peaks of β -Ga₂O₃ have a redshift, which indicates a corresponding tensile stress induced in the interface. In addition, the A_1 symmetry LO-phonon mode near 735 cm⁻¹ is not observed in the Raman spectrum of GaN/Ga₂O₃, which may be due to the strong coupling between the A_1 mode of the LO-phonon and the carrier plasmon (LO-phonon-plasmon coupling modes of electronic excitation and lattice vibrations).³⁸ Therefore, the LO-phonon-plasmon coupling causes the Raman peaks to be broader than those of a pure LO mode.³⁹ Furthermore, no Raman peaks related to the cubic GaN structure were observed. The Raman measurement therefore confirms the XRD finding that the GaN epilayer is a single-phase wurtzite, high-quality material.

In conclusion, we show a state-of-the-art GaN epilayer grown on (201) oriented β -Ga₂O₃ with high optical quality. The epitaxial relationship between the film and substrate showed low lattice mismatch (4.7%), which led to a low density of dislocations and other defects. The XRD and Raman

measurements showed a slight biaxial strain in GaN epilayers. The PL intensity was considerably higher than that of the GaN/Al₂O₃. This work shows that β -Ga₂O₃ is a promising substrate for the growth of high-quality GaN films. Under systematically optimized growth conditions, the growth of highly efficient LEDs and LDs based on III-nitrides should be possible and will allow new device designs such as vertically structured LEDs thanks to the increased transparency and high conductivity of bulk β -Ga₂O₃ substrates.

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